

Numerical Evaluation of Nonlinear Corrections to Nuclear Parton Evolution

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Inelastic Electron-Proton Scattering

Idea: Break protons up into their constituent particles
by scattering high-energy electrons off them

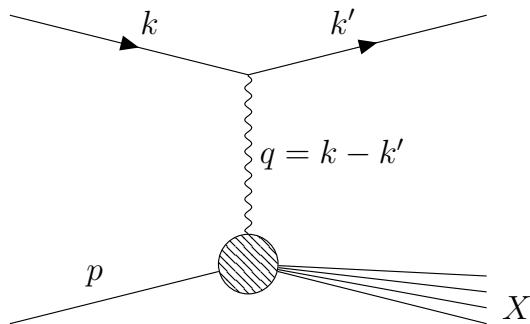
→ cross section:

$$d\sigma = \frac{1}{F} |\mathcal{M}|^2 dQ$$

$F \hat{=}$ incoming flux

$|\mathcal{M}|^2 \hat{=}$ $|i\rangle \rightarrow |f\rangle$ rate

$dQ \hat{=}$ $|f\rangle$ phase space



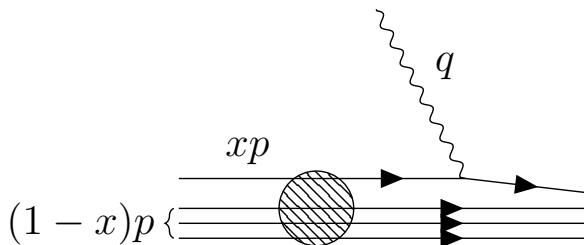
$$\text{Feynman rules} \implies d\sigma = \frac{\pi M}{(k \cdot p)} \frac{e^4}{q^4} L^{\mu\nu} W_{\mu\nu} \Pi^3 k' \quad \text{with}$$

$$W_{\mu\nu} = \frac{F_1}{M} \left(-g_{\mu\nu} + \frac{q_\mu q_\nu}{q^2} \right) + \frac{F_2}{\nu M^2} \left(p_\mu - \frac{p \cdot q}{q^2} q_\mu \right) \left(p_\nu - \frac{p \cdot q}{q^2} q_\nu \right)$$

Parton Model

Assumption: The proton consists of non-interacting quarks.

→ parton distribution functions (PDFs) $q_i(x)$



$q_i(x) \hat{=}$ probability quark i has momentum xp

If quarks are fermions, the cross section follows from QED.

⇒ Callan-Gross relation:

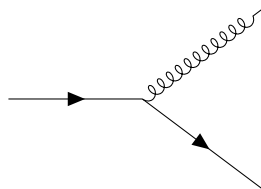
$$2xF_1(x) = F_2(x) = x \sum_{i=1}^{2N_F} e_i^2 q_i(x)$$

Quantum Chromodynamics

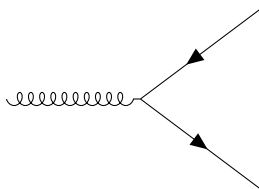
QCD is the theory of the interaction between quarks and gluons.

$$\mathcal{L}_{\text{QCD}} = \bar{\Psi}_f(i\not{D} - m_f)\Psi_f - \frac{1}{4}G_{\mu\nu}^a G_a^{\mu\nu}$$

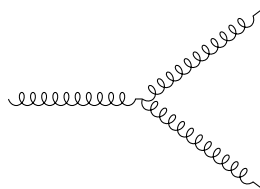
Processes affecting the PDFs:



q → g emission



g → q emission



g → g emission

→ PDFs now depend on the energy scale $Q^2 = -q^2$

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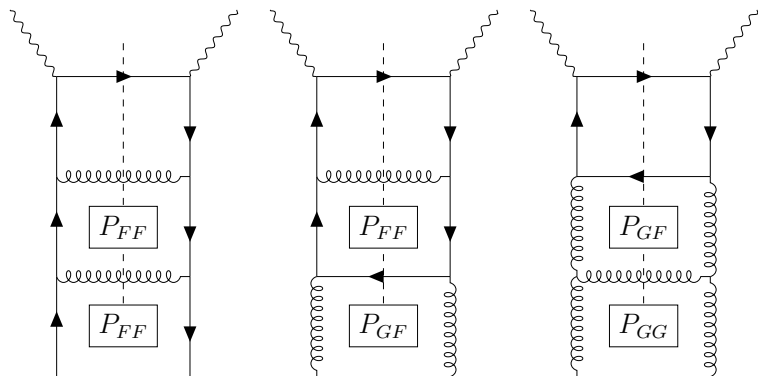
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DGLAP Equations

The DGLAP equations describe the Q^2 dependence of the PDFs.



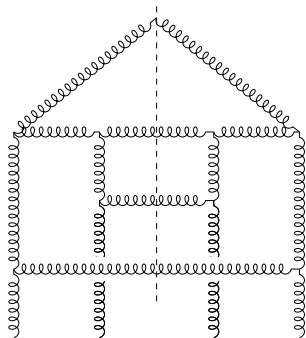
$$\frac{\partial}{\partial \ln Q^2} \begin{pmatrix} \Omega(x, Q^2) \\ G(x, Q^2) \end{pmatrix} = \frac{\alpha_s}{2\pi} \int_x^1 \frac{dz}{z^2} x \begin{pmatrix} P_{FF}(x/z) & P_{FG}(x/z) \\ P_{GF}(x/z) & P_{GG}(x/z) \end{pmatrix} \begin{pmatrix} \Omega(z, Q^2) \\ G(z, Q^2) \end{pmatrix}$$

with the momentum densities $\Omega = x\Sigma_f(q_f + \bar{q}_f)$, $G = xg$

GLR-MQ Equations

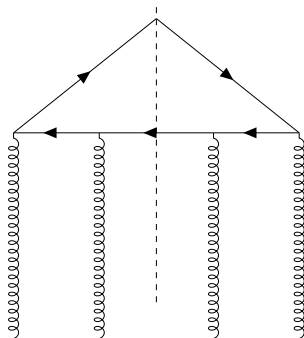
Taking gluon recombination into account modifies the evolution.

gluon-gluon recombination:



$$-\frac{81}{16} \frac{\alpha_s^2}{R^2 Q^2} \int_x^1 \frac{dz}{z} G^2(z, Q^2)$$

gluon-quark recombination:



$$-\frac{27}{160} \frac{\alpha_s^2}{R^2 Q^2} G^2(x, Q^2)$$

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Convolution Integrals

The evolution is done on a grid in x and Q^2 , with $D(x_c, Q_r^2) = D_{rc}$.

$$\int_{x_c}^1 dz P_{AB} \left(\frac{x_c}{z} \right) D(z) \xrightarrow[\text{Interpolation}]{\text{Linear}} \sum_{k=c}^n w_{AB}(x_k, x_c) D_{rk}, \text{ where}$$

$$w_{AB} \sim \int dz P_{AB}(z)$$

P_{AB} can be written as a perturbation series in α_s :

$$P_{AB}(x) = P_{AB}^{(0)}(x) + \frac{\alpha_s}{2\pi} P_{AB}^{(1)}(x) + \dots$$

Leading order contribution:

$$P_{GG}^{(0)}(z) = 2C_G \left[\frac{1}{1-z} + \frac{1}{z} - 2 + z - z^2 \right. \\ \left. - \frac{1}{z} \delta(1-z) \int_0^1 dy y \left\{ \frac{1}{1-y} + \frac{1}{y} - 2 + y - y^2 \right\} \right] - \delta(1-z) \frac{2}{3} T_{RN_F}$$

Convolution Integrals

Next-to-leading order contribution:

$$\begin{aligned} P_{GG}^{(1)} = & C_F T_R N_F \left[-16 + 8z + \frac{20}{3}z^2 + \frac{4}{3} \cdot \frac{1}{z} - (6 + 10z) \ln z - 2(1+z) \ln^2 z \right] \\ & + C_G T_R N_F \left[2 - 2z + \frac{26}{9} \left(z^2 - \frac{1}{z} \right) - \frac{4}{3}(1+z) \ln x - \frac{20}{9} \left(\frac{1}{1-z} + \frac{1}{z} - 2 + z - z^2 \right) \right] \\ & + C_G^2 \left[\frac{27}{2}(1-z) + \frac{67}{9} \left(z^2 - \frac{1}{z} \right) + \left(-\frac{25}{3} + \frac{11}{3}z - \frac{44}{3}z^2 \right) \ln z + 4(1+z) \ln^2 z \right. \\ & \left. + \left(\frac{67}{9} - \frac{\pi^2}{3} - 4 \ln z \ln(1-z) + \ln^2 z \right) \left(\frac{1}{1-z} + \frac{1}{z} - 2 + z - z^2 \right) + 2 \left(\frac{1}{1+z} - \frac{1}{z} - 2 - z - z^2 \right) \int_{z/(1+z)}^{1/(1+z)} \frac{dx}{x} \ln \left(\frac{1-x}{x} \right) \right] \\ & - \frac{1}{z} \delta(1-z) \int_0^1 dy y \left\{ C_F T_R N_F \left[-16 + 8y + \frac{20}{3}y^2 + \frac{4}{3} \cdot \frac{1}{y} - (6 + 10y) \ln y - 2(1+y) \ln^2 y \right] \right. \\ & + C_G T_R N_F \left[2 - 2y + \frac{26}{9} \left(y^2 - \frac{1}{y} \right) - \frac{4}{3}(1+y) \ln x - \frac{20}{9} \left(\frac{1}{1-y} + \frac{1}{y} - 2 + y - y^2 \right) \right] \\ & + C_G^2 \left[\frac{27}{2}(1-y) + \frac{67}{9} \left(y^2 - \frac{1}{y} \right) + \left(-\frac{25}{3} + \frac{11}{3}y - \frac{44}{3}y^2 \right) \ln y + 4(1+y) \ln^2 y \right. \\ & \left. + \left(\frac{67}{9} - \frac{\pi^2}{3} - 4 \ln y \ln(1-y) + \ln^2 y \right) \left(\frac{1}{1-y} + \frac{1}{y} - 2 + y - y^2 \right) + 2 \left(\frac{1}{1+y} - \frac{1}{y} - 2 - y - y^2 \right) \int_{y/(1+z)}^{1/(1+y)} \frac{dx}{x} \ln \left(\frac{1-x}{x} \right) \right] \left. \right\} \\ & - \delta(1-z) \int_0^1 dy y \left\{ C_F T_R N_F \left[4 - 9y + (-1 + 4y) \ln y + (-1 + 2y) \ln^2 y + 4 \ln(1-y) \right] \right. \\ & + \left(10 - \frac{2}{3}\pi^2 + 4 \ln y - 4 \ln y \ln(1-y) + 2 \ln^2 y - 4 \ln(1-y) + 2 \ln^2(1-y) \right) (y^2 + (1-y)^2) \left. \right\} \\ & + C_G T_R N_F \left[\frac{182}{9} + \frac{14}{9}y + \frac{40}{9} \cdot \frac{1}{y} + \left(\frac{136}{3}y - \frac{38}{3} \right) \ln y - 4 \ln(1-y) - (2 + 8y) \ln^2 y \right. \\ & \left. + \left(\frac{\pi^2}{3} - \frac{218}{9} + \frac{44}{3} \ln y + 4 \ln(1-y) - \ln^2 y - 2 \ln^2(1-y) \right) (y^2 + (1-y)^2) + 2 (y^2 + (1+y)^2) \int_{y/(1+y)}^{1/(1+y)} \frac{dx}{x} \ln \left(\frac{1-x}{x} \right) \right] \left. \right\} \end{aligned}$$

Evolution Process

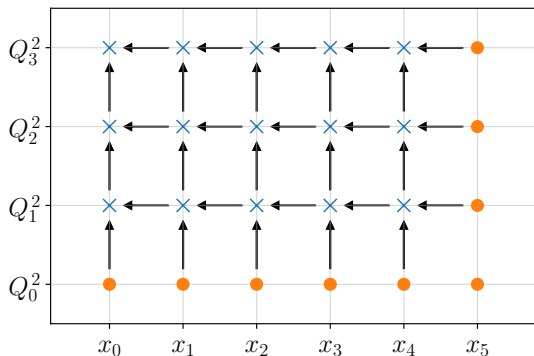
Evolution equations:

$$\Omega'_{rc} = W_{FF}\Omega_{rc} + W_{FG}G_{rc} + M_F - V_1G_{rc}^2$$

$$G'_{rc} = W_{GF}\Omega_{rc} + W_{GG}G_{rc} + M_G - (V_2G_{rc}^2 + V_3G_{rc} + N_G)$$

Two extra equations follow from a quadratic interpolation:

$$D_{rc} = D_{(r-1)c} + \frac{\Delta_r}{2} (D'_{(r-1)c} + D'_{rc})$$



Orange dots:
Starting values

At each point, the
four equations are
solved for D_{rc} & D'_{rc}

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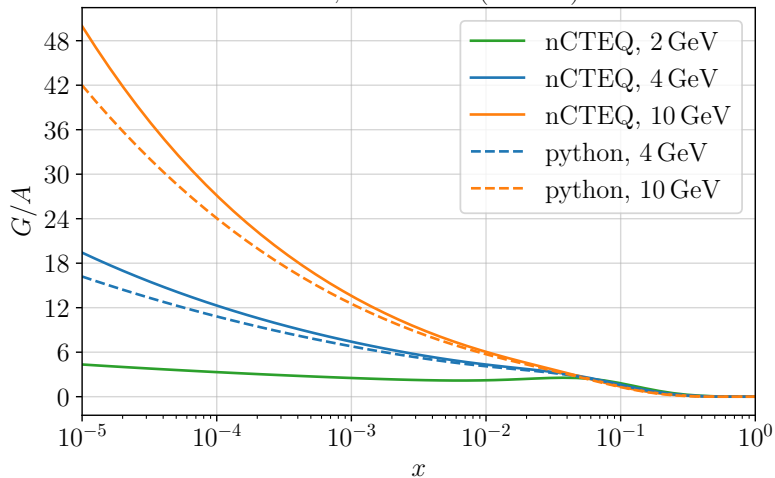
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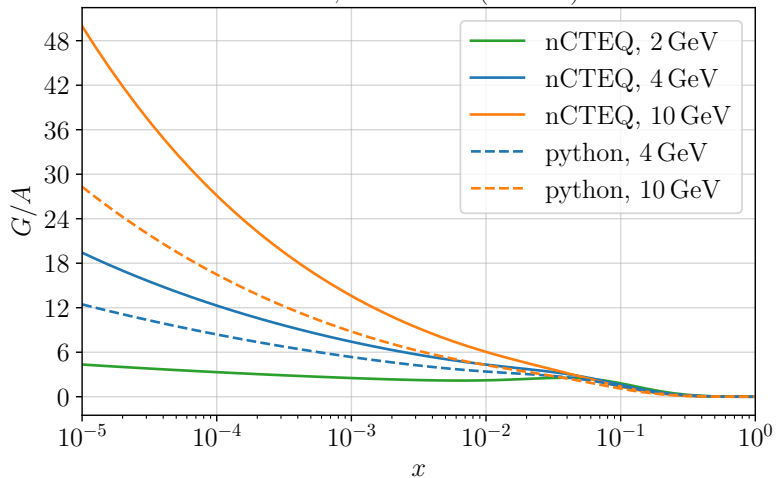
Interpretation

Gluon, nonlinear (Au197)



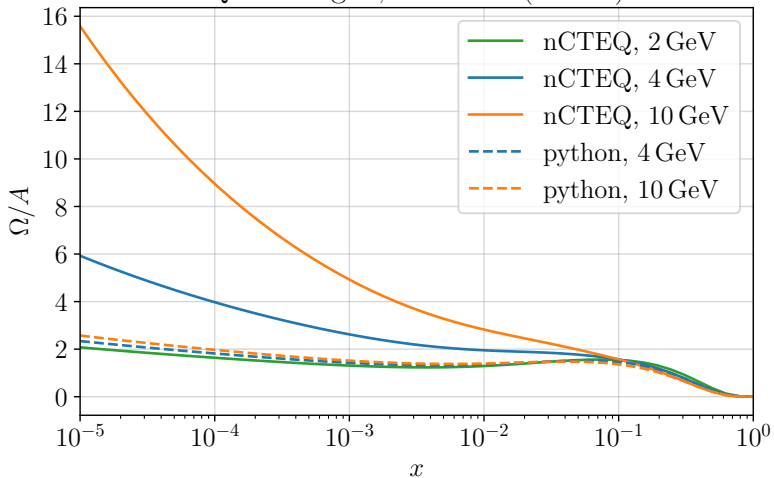
Parton Mixing

Gluon, nonlinear (Au197)

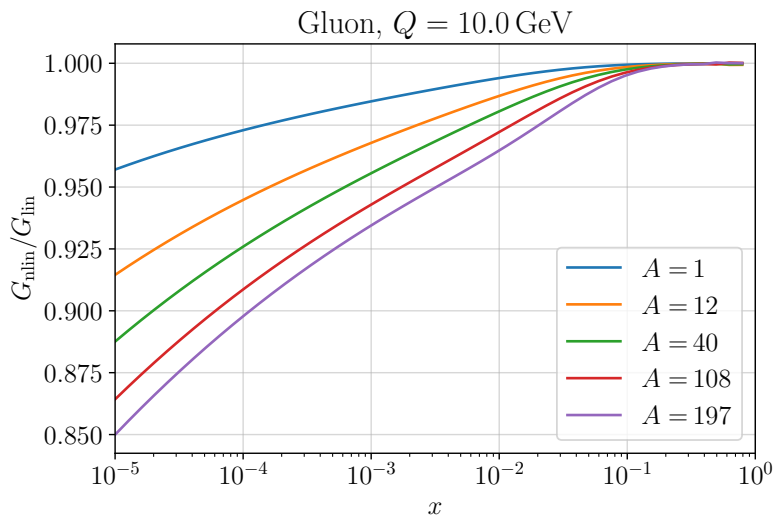


Parton Mixing

Quark singlet, nonlinear (Au197)



A Dependence



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